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Machine learning for detailed chemistry reduction in DNS of a syngas turbulent oxy-flame with side-wall effects

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Abstract

A novel chemistry reduction strategy based on convolutional neural networks (CNNs) is developed and applied to direct numerical simulation (DNS) of a turbulent non-premixed flame interacting with a cooled wall. The fuel syngas mixture is burning in pure oxygen. The training and the subsequent application of the CNN rely on the processing of two-dimensional (2D) images built from species mass fractions and temperature (CNN input), to predict the corresponding chemical sources at the center of the image (CNN output). This image-type treatment of chemistry is found to efficiently capture intermediate radicals species highly sensitive to the local flame topology. To reduce the CPU cost, a simplified 2D DNS database with detailed chemistry serves as reference and is used for training and testing the neural network. Comparisons are also made a posteriori against the same 2D DNS with a reduced chemical scheme specialized for syngas. Then, three-dimensional (3D) DNS are conducted either with CNN or the reduced chemistry for more a posteriori tests. The CNN reduced chemistry outperforms the reduced Arrhenius based mechanism in the prediction of radical species, such as monoatomic hydrogen, and also in terms of CPU cost.

Keywords: Convolutional neural network, Chemistry reduction, Direct numerical simulation, Syngas

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1. Introduction

The chemical mechanisms for hydrocarbon fuel combustion can involve up to hundreds of species and thousands of reactions [1]. Because of the very large number of degrees of freedom to be solved, such fully detailed mechanisms can hardly be introduced in three-dimensional (3D) numerical simulations of flames. In this context, the downsizing of chemical kinetics, still securing the prediction of given major thermochemical properties, is mandatory for virtual prototyping of real combustion systems, particularly those involving heavy hydrocarbon liquid or solid fuels [2].

Numerous successful strategies have been discussed in the literature to reduce detailed chemical kinetics. They are based mainly on fuel lumping, time-scale and principal component analysis or direct numerical treatment of the chemical responses from graph analysis or genetic algorithms, sometimes associated with tabulation methods [3–17].

Learning from a database to then reproduce the chemical response has received extensive attention in tabulation of combustion chemistry. Artificial neural networks (ANNs) were found accurate, and both CPU and memory efficient [18, 19]. Along similar lines, convolutional neural networks (CNNs), which were originally developed for analyzing visual representations [20], have been introduced in turbulent combustion modeling for direct deconvolution of the filtered progress variable [21], for modeling the unresolved flame surface wrinkling [22] and for the direct reconstruction of unresolved sources and fluxes from mesh-resolved quantities in large-eddy simulation (LES) [23]. Compared to ANNs, CNNs reduce the number of connections per layer to stack more of them efficiently, thus increasing the depth of the neural network.

In the present study, a novel chemistry reduction approach based on CNNs is proposed and also coupled with a three-dimensional direct numerical simulation (DNS) for validation. A turbulent non-premixed oxy-flame interacting with a cooled wall is considered. This configuration covers a wide range of combustion and turbulence conditions in terms of equivalence ratio, progress of reaction and enthalpy including thermal and viscous boundary layers effects, making it a relatively generic test case. The syngas fuel composition is representative of a CO₂ free blast-furnace exhaust flux, as found in a recycling combustion process. The input of the CNN are two-dimensional (2D) images of the species mass fractions and temperature, while the output are the corresponding chemical burning rates at the center point of the image. The training is performed from

a simplified two-dimensional planar DNS, with detailed chemistry and input images built from a subset of the full species vector. This subset of 11 species is chosen to match the major and minor species retained in a reduced chemical mechanism optimized for syngas. Once the CNN trained, this reduced set of species becomes the input to get the corresponding chemical rates. Then, only this reduced set of species needs to be transported with the flow, saving much CPU time and reducing memory requirement.

The objective is to test the feasibility of the novel method, which should be applicable to any reference detailed chemical mechanism. The generic GRI-3.0 detailed mechanism [24] without NO_x (217 elementary reactions and 35 chemical species) is used in this first attempt. Once trained for reproducing the chemical sources of a reduced set of 11 species, the CNN for chemistry reduction is tested a priori against the detailed chemistry DNS. Then, the obtained neural network is fully coupled with the flow solution for two- and three-dimensional DNS of the non-premixed flame-wall interaction. The two-dimensional cases are with the detailed chemistry, the reduced chemical mechanism and the CNN chemistry. The three-dimensional cases are with the CNN and the reduced chemistry.

2. Turbulent flame configuration and numerics

A turbulent non-premixed oxy-flame flame under a pressure of 341.3 kPa and featuring side-wall effects is simulated (Fig. 1). A representative syngas composition is injected through the upper inlet (width of 2 mm) at 1223 K and bulk velocity of 50 m/s, with H₂/0.0085, CO/0.7852, CO₂/0.0514, N₂/0.1549 in mass fraction. Below the splitter plate of thickness 0.2 mm, the lower inlet (width of 10 mm) supplies pure oxygen at 298 K and 100 m/s. A preliminary simulation of an oxygen turbulent periodical 3D channel flow at the oxygen bulk velocity serves as a database to impose a proper turbulent inlet to the oxygen stream. A laminar channel flow profile is imposed for the fuel stream, between the wall and the splitter plate. A non-slip wall is placed at the top of the domain while a symmetrical boundary condition is used at the bottom. Periodic boundary conditions are employed in the spanwise z direction (Fig. 1) and a convective boundary condition is used for the outflow. The wall temperature $T_w(\underline{x}, t)$ is non-uniform and computed assuming that the heat flux locally reaching the wall is evacuated on its other side by an efficient water cooling with a convective heat transfer $h = 3300 \text{ W} \cdot \text{m}^{-2} \cdot \text{K}^{-1}$, with a water temperature $T_o = 353 \text{ K}$ (i.e., $h(T_w - T_o) = \lambda(\partial T/\partial n_w)_{gas}$, n_w is the

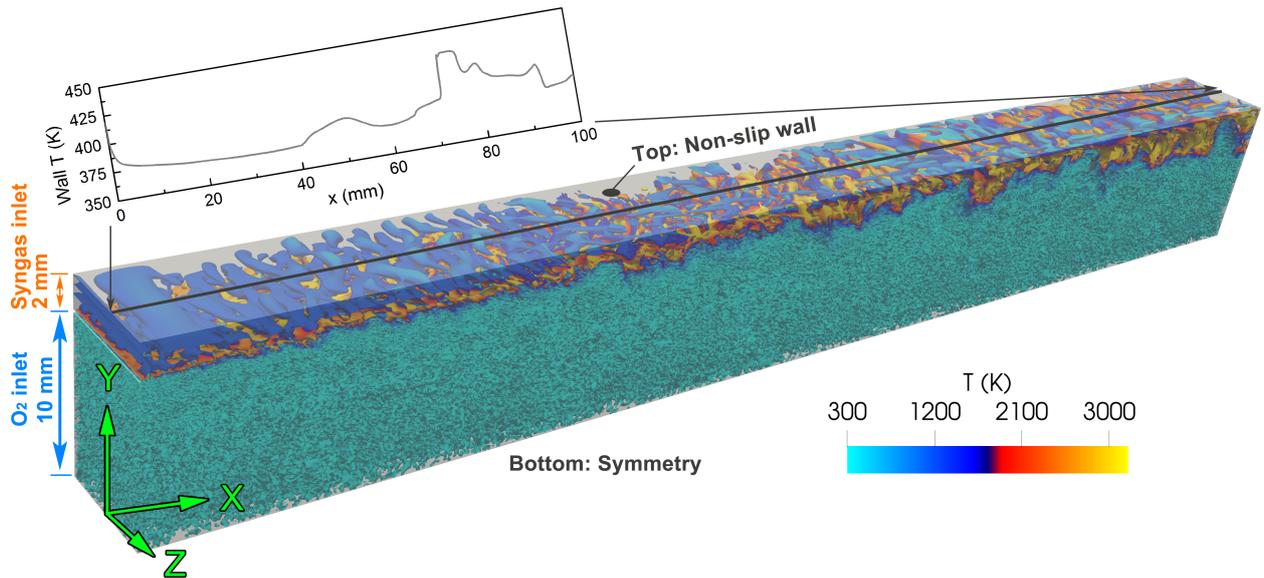


Figure 1: Instantaneous iso-surface of Q criterion [25] predicted by CNN-DNS: $Q = 10^7 \text{ s}^{-2}$, colored by gas temperature. Top graph: distribution of instantaneous wall temperature along the centreline.

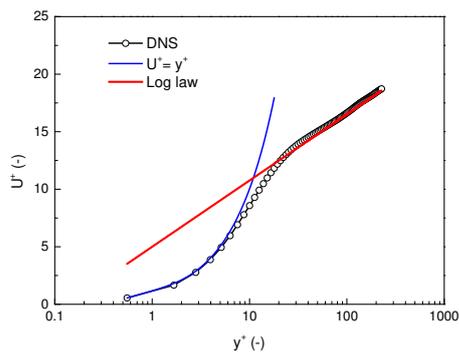


Figure 2: Velocity distribution in wall units for 3D non-reactive case (mesh verification).

wall normal direction). The physical dimensions of the computational domain composed of 66.44M cells are $L_x = 100 \text{ mm}$, $L_y = 12 \text{ mm}$ and $L_z = 7.68 \text{ mm}$. The mesh is uniform in x (streamwise) and z (spanwise, not considered in 2D) directions, with the grid spacing $h_x = h_z = 60 \mu\text{m}$. In the y (vertical) direction, the mesh keeps a fine resolution of $20 \mu\text{m}$ for the upper region close to the wall ($y > 10 \text{ mm}$) where the flame is located.¹ The mesh gets coarser as y decreases from 10 to 0 mm to reach a grid spacing of $100 \mu\text{m}$ at the bottom of the computational domain. The first point off the

¹ Simulations of one-dimensional canonical strained non-premixed flames were conducted to determine the grid resolution for the present free stream conditions.

top wall is at $y^+ = 0.55$ where the superscript + denotes the usual non-dimensionalization by the viscous length scales. There are 9 grid points within $y^+ = 10$ to resolve the viscous layer near the wall [26]. It was verified from a non-reactive simulation that the structure of the wall boundary layer agrees with what is expected (Fig. 2). The time step is fixed at $0.3 \mu\text{s}$ ($\text{CFL} \approx 0.8$) with sub-iteration in a stiff-chemistry ODE solver.

The governing equations for the gas mixture are solved with a low-Mach number in-house code [27–31]. The conservation equations for mass, momentum, species and temperature of the gas phase are those of [31], here without the source terms due to coal particles. The flow viscosity and the thermal conductivity of the gas mixtures are computed from Wilke’s type mixing rules. The Lewis number is assumed to be unity for all the species.

The numerics is based on an approach previously used for both DNS and LES [32, 33]. The time advancement relies on a second-order Crank-Nicolson scheme. A bounded quadratic upstream interpolation for convective kinematics (BQUICK) scheme is used for the scalar advection terms in the species and temperature equations, while a second-order central difference scheme is applied to the scalar diffusion terms in the species and temperature equations and all terms in the momentum equation. An alternating direction implicit (ADI) method has been employed, in which semi-implicit tridiagonal/pentadiagonal equations are solved

Table 1: Structure of the sequential CNNs used (see [34] for details on functions).

Layer type	Activation function	Output Shape
Input	–	(None, 3, 3, 11)
Conv2D	ReLU	(None, 3, 3, 64)
MaxPooling2D	–	(None, 2, 2, 64)
Conv 2D	ReLU	(None, 2, 2, 256)
MaxPooling2D	–	(None, 1, 1, 256)
Flatten	–	(None, 256)
Dense	ReLU	(None, 256)
Dense	ReLU	(None, 64)
Dense	ReLU	(None, 16)
Dense (Output)	–	(None, 11)

separately for each direction.

3. CNN chemistry reduction methodology

3.1. Training phase

Figure 3 sketches the overall CNN training process from the 2D database, which is applied to the species O, O₂, H, H₂, OH, HO₂, H₂O, H₂O₂, CO and CO₂ (NO_x being not included, no source is computed for N₂). Because the syngas is preheated to 1223 K, rapid ignition occurs after mixing with pure oxygen and a non-premixed flame develops in the computational domain (Fig. 1). After the flame is established, the simulation proceeds for 2 flow-through times (2 ms) based on the bulk velocity of the O₂ stream.

The training database is built saving every 50 μs the 0.52 million 2D DNS fields of species mass fractions and temperature together with their sources, leading to 40 snapshots each containing 520,000 data. To augment the database up to 80 snapshots, the input data are perturbed applying 2% of uncorrelated random noise. Then, over the 80 snapshots, one of every 5 is selected to form the subsequent testing database (16 snapshots), while the left 64 snapshots constitute the training database.

In order to train the CNN with the two-dimensional simulations, a 2D test-box is constructed around every i -th DNS node with its neighboring 8 nodes (Fig. 3). Selecting the 8 neighboring nodes with the target node at the center, to form a 3x3 2D image, appears as the most efficient procedure with the smallest size of CNN input. This test-box is used to build the i -th image composed of 9 points from the 2D distributions of the vector $\underline{\phi}(x_\ell, t) = (Y_1(x_\ell, t), \dots, Y_N(x_\ell, t), T(x_\ell, t))$

for $\ell = 1, \dots, 9$, which contains information on the $N = 10$ species mass fractions and temperature (N₂ not included). The “target” (or “label”) of each image is the vector of the reaction rates at the \underline{x} -centred i -th point, $\underline{\omega}[i] = (\underline{\phi}(\underline{x}, t + \delta t) - \underline{\phi}^*(\underline{x}, t))/\delta t$, with $\delta t = 0.3\mu\text{s}$ and where $\underline{\phi}^*$ denotes the solution before the ODE stiff-chemistry solver is applied. Casting the “targets” (i.e., the sources) in this form presents many advantages. Indeed, the output of the CNN can be used either as a chemical source $\underline{\omega}[i]$ for integration, or directly to read from the CNN $\underline{\phi}(t + \delta t) = \underline{\omega}[i] \times 0.3 \cdot 10^{-6} + \underline{\phi}^*(t)$ in a flow solver relying on splitting for chemistry integration where $\underline{\phi}^*(t)$ is known. In the latter option, sub-iterations may be needed for $\delta t > 0.3\mu\text{s}$. In our simulations, the time step stays the same between training and usage of the CNN, therefore no linear approximation is actually performed. However, care should be taken that splitting method for chemistry integration can decrease the overall order of accuracy of the solver.

To secure the atomic mass balance over a reduced set of species whose sources ω_i are learned from a detailed mechanism, the following procedure is adopted. (In practice, over the computational domain, the total default mass to be corrected represents only 0.016% of the total mass.) First, the atomic sources mass default are computed for carbon, hydrogen and oxygen. Let us denote $\alpha_{A,i}$ the number of atom A (i.e., $A \equiv \text{C, H or O}$) in the i -th species. The mass imbalance in A source may be written $\Delta\omega_A = \sum_{i=1}^N [(\alpha_{A,i}W_A)/W_i]\omega_i$, where W_A and W_i are molar weights. Second, sources of carbon containing species of the syngas (CO and CO₂) and hydrogen containing species (H₂, H, H₂O, OH, HO₂ and H₂O₂) are corrected adding $\Delta\dot{w}_i = -(Y_i/Y_A)\Delta\omega_A$, with the atom mass fraction $Y_A = \sum_{i=1}^N \alpha_{A,i}(W_A/W_i)Y_i$. For oxygen atom, the correction is applied only to O₂ and with the already corrected sources of C and H containing species. Then, atomic budgets are perfectly preserved by the chemical sources learned by the CNN.

46% of the images have a central node temperature below 320 K, these points with negligible sources are dropped out from the database. This leads to a final training database of 14.6 GB and a testing database of 3.7 GB. All data contributing to the “images” and to the “targets” are normalized by their respective maximum absolute value over all the database, these 22 maximum levels are stored (10 species mass fractions, temperature and sources – N₂ not involved). The complete structure of the sequential CNNs is summarized in Table 1. The CNN architecture is inspired from [23], with two additional dense layers to secure the performance of the regression task required here. In total, the neural network

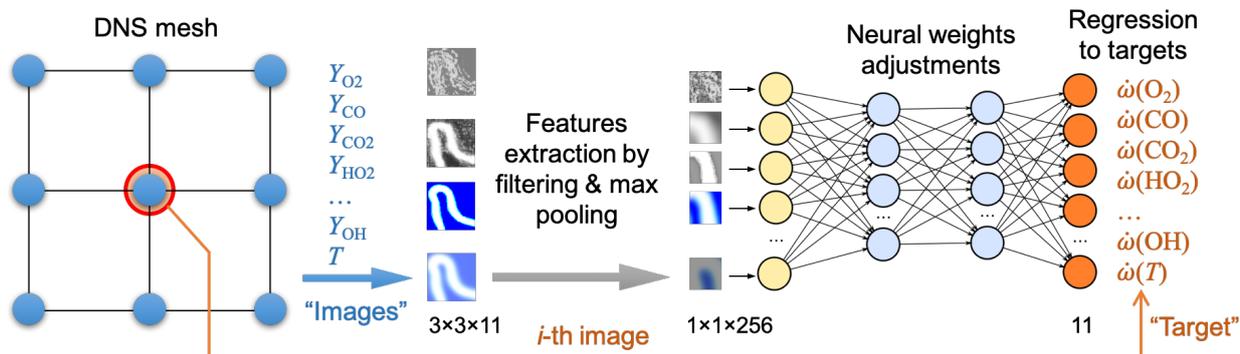


Figure 3: CNN training from 2D database.

contains 237,579 weights, which need to be adjusted.

To improve the precision of the CNNs for small values of the chemical sources, a second network with exactly the same structure is trained for a subset of the database defined by $Y_{O_2} > 0.9$ at the central node (i.e., where overall the burning rates stay very low). The subset-database is normalized based on its maximum absolute values. In the following, the networks are referred to as the main CNN and the second CNN (subset for lower burning rates).

The training of the CNNs is performed using the TensorFlow Python library with GPU support (www.tensorflow.org). A series of convolution/sampling operations are performed iteratively, in which the neural weights are adjusted until a satisfying minimal error is achieved between the “target” sources and the ones predicted by the CNNs. The Adam optimizer [35] is employed with a mean-squared-error (MSE) loss function comparing the predictions to the target. A training step is performed with a batch of 10240 “images”, leading to 1740 steps per “epoch”/iteration for the main CNN. The learning rate [34] is initially set to 0.001 to decay versus the steps (n) according to $l_r = 0.001/(1.0 + 0.0001 \cdot n)$. The network converges in approximately 500 epochs, requiring for the main CNN 15 hours on an Nvidia GeForce GTX 1080 Ti GPU.

3.2. Testing phase

The performance of the machine learning reduced chemistry is first examined in a-priori manner by comparing chemical sources from CNN and GRI-3.0. The part of the database unseen during training is probed. The mean-square error between the predictions and the targets is 1.5×10^{-7} for the main CNN and 4.5×10^{-8} for the second CNN. Figure 4 displays the predicted

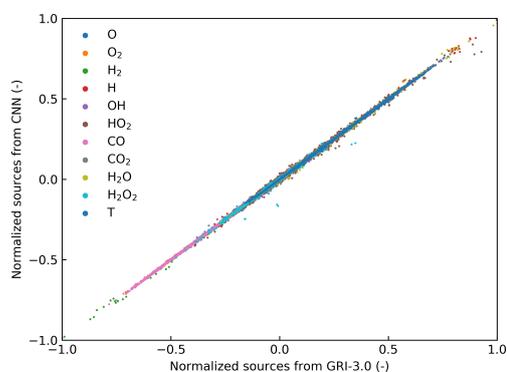


Figure 4: Comparison of the main CNN prediction on normalized sources of the 10 species and temperature against the GRI-3.0 target values for the testing database unseen during training.

chemical sources against their reference values for the main CNN, confirming the accuracy of the approach. The second CNN achieves similar performance in this a-priori validation test (not shown for brevity).

4. CNN-DNS full coupling

The CNN reduced chemistry is now employed in the simulation of the syngas non-premixed flame (only CPU is used in this part with CNNs exported in C language, see supplementary material). Simulations are then performed with detailed chemistry (2D), CNN chemistry (2D and 3D) and a reduced chemistry (2D and 3D) given in Table 2. This reduced mechanism was derived from GRI-3.0 for the present syngas operating condition including wall heat loss effects, with the automated approach reported in detail in [13, 36].²

²This method combines a stochastic probing of the detailed chemistry response with directed relation graphs analysis with error prop-

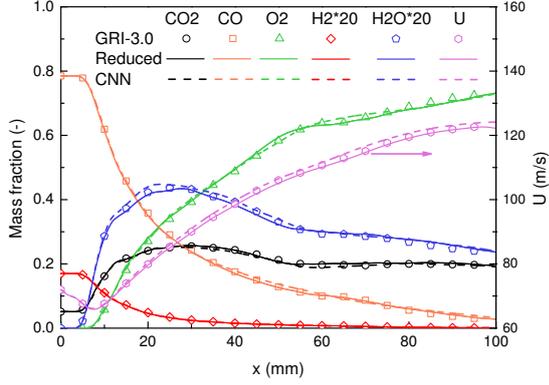


Figure 5: Averaged distributions of streamwise velocity and major species mass fractions: CO_2 , CO , O_2 , H_2 and H_2O along the planar fuel-jet centreline ($y = 11$ mm). Symbols: GRI-3.0. Solid line: Reduced mechanism (Table 2). Dashed line: CNN chemistry, 2D case.

Table 2: Reduced chemistry for syngas oxy-flame. Units are mol, s, cm^3 , cal and K. The Chaperon efficiencies of the GRI-3.0 mechanism are preserved for both three-body and fall-off reactions (see also supplementary material).

n^o	Reaction	A	β	E_a
1	$2\text{O} + \text{M} \rightleftharpoons \text{O}_2 + \text{M}$	$1.20\text{e}+17$	-1.00	0
2	$\text{O} + \text{H}_2 \rightleftharpoons \text{H} + \text{OH}$	$3.87\text{e}+04$	2.70	6260
3	$\text{O} + \text{HO}_2 \rightleftharpoons \text{OH} + \text{O}_2$	$2.00\text{e}+13$	0	0
4	$\text{O} + \text{CO} (+\text{M}) \rightleftharpoons \text{CO}_2 (+\text{M})$	$1.80\text{e}+10$	0	2385
5	$\text{O}_2 + \text{CO} \rightleftharpoons \text{O} + \text{CO}_2$	$2.50\text{e}+12$	0	47800
6	$\text{H} + \text{O}_2 + \text{M} \rightleftharpoons \text{HO}_2 + \text{M}$	$2.80\text{e}+18$	-0.86	0
7	$\text{H} + \text{O}_2 + \text{N}_2 \rightleftharpoons \text{HO}_2 + \text{N}_2$	$2.60\text{e}+19$	-1.24	0
8	$\text{H} + 2\text{O}_2 \rightleftharpoons \text{HO}_2 + \text{O}_2$	$2.08\text{e}+19$	-1.24	0
9	$\text{H} + \text{O}_2 + \text{H}_2\text{O} \rightleftharpoons \text{HO}_2 + \text{H}_2\text{O}$	$1.13\text{e}+19$	-0.76	0
10	$\text{H} + \text{O}_2 \rightleftharpoons \text{O} + \text{OH}$	$2.65\text{e}+16$	-0.67	17041
11	$\text{H} + \text{OH} + \text{M} \rightleftharpoons \text{H}_2\text{O} + \text{M}$	$2.20\text{e}+22$	-2.00	0
12	$\text{H} + \text{HO}_2 \rightleftharpoons \text{O}_2 + \text{H}_2$	$4.48\text{e}+13$	0	1068
13	$\text{H} + \text{HO}_2 \rightleftharpoons 2\text{OH}$	$8.40\text{e}+13$	0	635
14	$\text{OH} + \text{H}_2 \rightleftharpoons \text{H} + \text{H}_2\text{O}$	$2.16\text{e}+08$	1.51	3430
15	$2\text{OH} (+\text{M}) \rightleftharpoons \text{H}_2\text{O}_2 (+\text{M})$	$7.40\text{e}+13$	-0.37	0
16	$\text{OH} + \text{HO}_2 \rightleftharpoons \text{O}_2 + \text{H}_2\text{O}$	$1.45\text{e}+13$	0	-500
17	$\text{OH} + \text{H}_2\text{O}_2 \rightleftharpoons \text{HO}_2 + \text{H}_2\text{O}$	$2.00\text{e}+12$	0	427
18	$\text{OH} + \text{H}_2\text{O}_2 \rightleftharpoons \text{HO}_2 + \text{H}_2\text{O}$	$1.70\text{e}+18$	0	29410
19	$\text{OH} + \text{CO} \rightleftharpoons \text{H} + \text{CO}_2$	$4.76\text{e}+07$	1.23	70
20	$2\text{HO}_2 \rightleftharpoons \text{O}_2 + \text{H}_2\text{O}_2$	$1.30\text{e}+11$	0	-1630
21	$2\text{HO}_2 \rightleftharpoons \text{O}_2 + \text{H}_2\text{O}_2$	$4.20\text{e}+14$	0	12000
22	$\text{HO}_2 + \text{CO} \rightleftharpoons \text{OH} + \text{CO}_2$	$1.50\text{e}+14$	0	23600
23	$\text{OH} + \text{HO}_2 \rightleftharpoons \text{O}_2 + \text{H}_2\text{O}$	$5.00\text{e}+15$	0	17330

- **2D simulations:** The results of CNN-DNS are compared to that of the detailed GRI-3.0 and the 11-species reduced mechanisms of Table 2 in Figs. 5 and 6 (streamwise and spanwise distributions). Statistical means are computed after the flow is established (10 flow through times). The averaged streamwise velocity distribution along the planar fuel jet centerline decreases until $x \approx 8$ mm (Fig. 5), location where the fuel ignites and

agitation [11, 37] and a genetic algorithm [38]. The reaction rates of the reduced mechanism are optimized against the response of the detailed mechanism in non-adiabatic non-premixed canonical problems involving turbulent micro-mixing as in [36].

the mass fractions of CO and H_2 decrease. The non-premixed turbulent flame scalar structure is as expected, the syngas fuel composed of CO and H_2 is mainly transformed into CO_2 and H_2O (Fig. 5), with a slowing down of chemistry close to the cooled wall. Because of the overall fuel-lean character of the system, the mass fraction of O_2 increases along the centreline.

According to Figs. 5 and 6, the coupled CNN-DNS reproduces very well the scalar structure of the non-premixed planar jet flame, with a close agreement against the same DNS with the detailed GRI-3.0 mechanism. Interestingly, for the same set of species transported with the flow, the reduced mechanism tends to over-predicts the mass fraction of H radical while the CNN-DNS keeps a close agreement with the reference detailed mechanism (Fig. 6). The prediction of the radical H chemical source may actually benefit from the image-type treatment by CNN, in which the 2D distribution of species is used as input. The averaged departure of wall heat flux between detailed and CNN (respectively reduced) chemistry is 3% (respectively -1%). With the CNN, the standard deviation of the departure represents 6.5% of the mean heat flux.

- **3D simulations:** The statistics of the 3D run with CNN and the reduced mechanism are initialized from scratch and collected from 1 flow through time, with additional averaging over the 128 spanwise planes, after the initial condition has been evacuated from the computational domain and the turbulence is fully developed. Thereby the memory of the initial condition cannot affect the results. Even though it is sufficient for training, the two-dimensional channel turbulence lacks vortex stretching and the three-dimensional turbulent flow cascade differs in 3D, modifying the turbulence/chemistry interaction and the scalar averages, as seen comparing Figs. 6 and 7. The predictions of the CNN-DNS still match those by the reduced mechanism DNS, but for H mass fraction, the reduced mechanism predicting again a higher level than the CNN. Based on above comparisons against GRI-3.0 in 2D, it is legitimate to assume that the scalar flame structure obtained with the CNN chemistry is actually closer to the one that would be seen in 3D with the detailed mechanism (out of reach with actual computing power). The averaged departure of wall heat flux between CNN and reduced chemistry is 14%, with a standard deviation of 13.8% of the mean flux.

The comparison of the computational costs due to chemistry for advancing the thermochemical variables, thus excluding the flow transport part, shows that CNN-chemistry is 11 times faster than the GRI-3.0 detailed mechanism, and 1.35 times faster than the 11-species reduced mechanism, the latter mainly because the Ar-

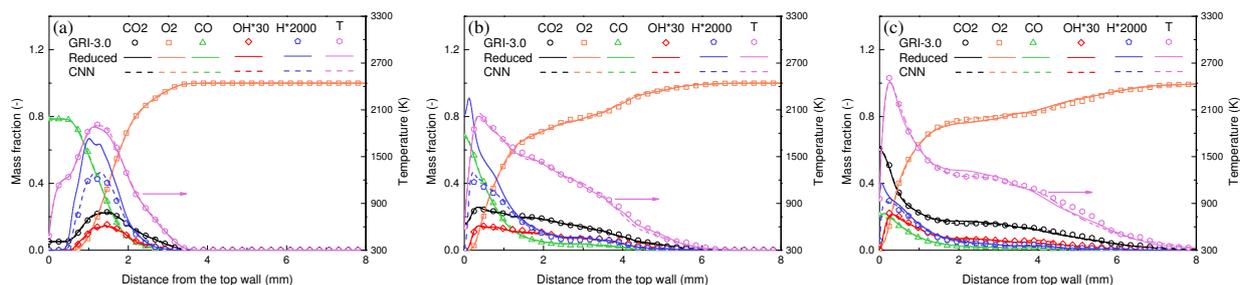


Figure 6: Spanwise averaged distributions of CO_2 , O_2 , CO , OH , H mass fractions and temperature. Symbols: GRI-3.0. Solid line: Reduced mechanism (Table 2). Dashed line: CNN chemistry. (a): 15 mm from the nozzle inlet. (b): 50 mm from the nozzle inlet. (c): 80 mm from the nozzle inlet. 2D case.

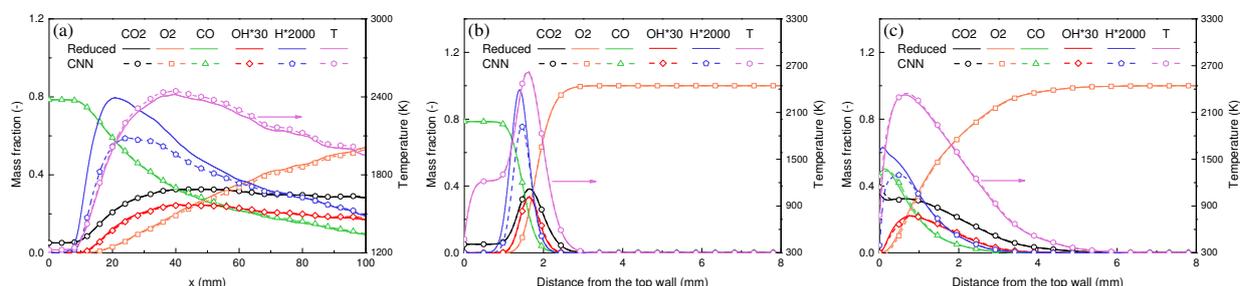


Figure 7: Averaged distributions of CO_2 , O_2 , CO , OH , H mass fractions and temperature. Solid line: Reduced mechanism (Table 2). Line with symbols: CNN chemistry. (a): BFG jet centerline (streamwise). (b): 15 mm from the nozzle inlet (spanwise). (c): 80 mm from the nozzle inlet (spanwise). 3D case.

Arrhenius rates calculation is avoided.

5. Summary

A new paradigm grounded on machine learning is discussed to introduce complex chemistry in turbulent flame simulation and implemented in three-dimensional direct numerical simulation of a syngas non-premixed flame/wall interaction. The training of a convolutional neural network (CNN) is performed on a simplified two-dimensional simulation with a reference detailed chemistry. The obtained reduced chemistry is then applied to a three-dimensional turbulent flame simulation.

The comparisons against the detailed chemistry response, a reduced mechanism and the CNN confirm the validity of the approach, which allows for conveniently developing reduced chemistry specialized for given operating conditions from a generic detailed scheme and even save CPU time compared to using reduced mechanisms based on Arrhenius rates.

One major drawback of the presented approach lies in the need to simulate a simplified two-dimensional version of the target three-dimensional flow geometry to train the CNN. However, in future works, generic training procedures will certainly emerge, so that actually

performing a flow simulation will not be mandatory to optimize the neural networks. Along these lines, the use of non-premixed stochastic micro-mixing model problems appears as a promising alternative to perform CNN training in a pre-processing mode. Because it has already been successfully used as a canonical problem, prior to any flow simulation, in the context of automated chemistry reduction [13].

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